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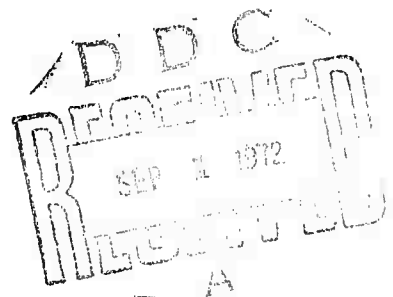
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13. ABSTRACT

The emission spectrum of the active nitrogen-nickel carbonyl reaction flame between 2300A and 7600A has been measured in a discharge/flow system at pressures between 0.13 and 27 torr. Relative populations and relative rates of formation of excited states of atomic nickel have been calculated, and various effects of reaction conditions and gas composition have been investigated. To the extent testable, the results are consistent with a mechanism in which atomic nickel produced in the reaction is excited by collision with metastable $N_2(A^3\Sigma_u^+)$ molecules in various vibrational levels produced in active nitrogen. The pressure dependence of the atomic nickel spectrum results mainly from changes in the vibrational distribution of the $N_2(A)$ molecules.

KEY WORDS

Atomic nickel excitation
 Active nitrogen reaction
 Quantitative emission spectroscopy
 Nickel carbonyl

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Final Technical Report

EXCITATION OF ATOMIC NICKEL IN THE REACTION
BETWEEN NICKEL CARBONYL AND ACTIVE NITROGEN

by

E. C. Shane and W. R. Brennen

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INTRODUCTION

Since the observation of Strutt¹ that the line spectra of various metals are developed when the metals are exposed to active nitrogen, several studies^{2,3,4,5} of excitation of metal atoms in active nitrogen have been made. Brennen and Kistiakowsky⁵ found that Ni(CO)_4 , Fe(CO)_5 , Cr(CO)_6 , W(CO)_6 , $\text{Mn}_2(\text{CO})_{10}$, and Co(NO)(CO)_3 reacted rapidly with active nitrogen, forming metal atoms, and producing flames due to metal atom emission spectra. We have looked in greater detail at the Ni(CO)_4 -active nitrogen reaction. Emission spectra were observed between 0.13 and 27 torr and found to be pressure dependent. A simple excitation mechanism is presented and the effect of quenching and the applicability of the Franck-Condon principle to the energy transfer reaction are discussed. The effect of reactant ratio, observation time, and added gases on the atomic flame are evaluated.

EXPERIMENTAL SECTION

The Ni(CO)_4 -active nitrogen reaction was studied in a 25mm o.d. pyrex flow tube fitted with a fused silica observation window mounted parallel to the axis of the flow tube as shown in Fig. (1). Gas inlet jets were located immediately before the observation window and 19 cm upstream. For one experiment an inlet jet was located 5.3 cm upstream from the observation instead of right at it. Downstream from the flow tube were a 15 mm bore, right angle throttling stopcock, two U-traps, one of which was detachable, a 10 mm bore right angle stopcock, a 2-liter ballast bulb, and a mechanical pump. While Ni(CO)_4 was flowing both U-traps were maintained at liquid nitrogen temperature. Light emitted by the reaction flame passed through the fused silica window and was directed by two reflecting mirrors into a 3/4 meter

Czerny-Turner, Jarrell-Ash scanning monochromator located about 34 cm distant. For wavelengths between 4000 and 7600 Å a Corning 3060 filter was inserted just before the entrance slit of the monochromator. Primarily photoelectric measurements were made using a dry ice-cooled EMI-9558 photomultiplier tube with magnetic defocusing operated at 860V from a Northeast Model RE-1602 stabilized high-voltage power supply. The output of the photomultiplier tube was amplified by a Keithley Model 610-BR electrometer and displayed on a Leeds and Northrup Speedomax-G strip chart recorder. Typical monochromator slit widths were 152-200 μ , giving a band pass of about 1.5 to 2.0 Å. A few photographic measurements were made using 50 μ slits and Kodak 103-F photographic plates.

The relative sensitivity between 2500 and 7500 Å of the 0.75 m monochromator and EMI-9558QA photomultiplier tube was calibrated with a standard tungsten-iodine lamp from the Eppley Corporation. Over a one-year period no change in the relative sensitivity was detected. A plot of the sensitivity, defined as the photomultiplier tube output in amps at a given wavelength divided by the standard lamp output in quanta/sec, against wavelength is given in Fig. (3). A correction for a slight change in reciprocal linear dispersion over the wavelength range is included.

Nitrogen atoms were formed in either an electrodeless, microwave discharge in a fused silica tube located 70 cm upstream from the flow tube or a pulsed discharge tube located 140 cm upstream. The microwave discharge was powered by a 2450 MHz, 100 W Raytheon Model PGM-10 microwave power generator coupled to a Type 5⁶, 1/4 wave cavity supplied by Ophthos Instruments. The pulsed discharge was maintained in a water-cooled, pyrex discharge tube, Fig. (2a), fitted with two stainless steel hollow electrodes. The condensed discharge circuit, Fig. (2b), employed a Sorenson High Voltage Supply Model 1006-100

with a maximum output of 6000V and 100mA. The power supply charged a Tobe Deutschmann 1 μ F, 10,000V capacitor. The rate of firing was regulated by a spark gap with two opposing brass rods with hemispherical ends separated by a rotating brass conductor. The entire spark gap assembly was enclosed in a plexiglass box which was flushed with nitrogen. The rate of rotation of the center conductor was regulated by a series motor powered by a Variac. The rpm of the motor and thus the pulse rate was easily regulated with the Variac. Typically the pulsed discharge was run at 5000V with a pulse rate of 8-15 pulses/sec. After several hours of running the nitrogen atom concentration began to fall. Apparently metal which sputtered from the electrodes deposited on the exit tube where it enhanced wall recombination of the nitrogen atoms. Rinsing with a 10% aqueous HF solution removed any deposited metal.

The nitrogen atom concentration was determined by NO titration using a visual endpoint. The flow rate of NO, $\text{Ni}(\text{CO})_4$, and added gases such as CO and C_2H_2 were measured by timing a pressure drop in a calibrated volume⁷ using Dow-Corning DC-704 silicone oil ($\rho=1.063$ g/cc at 25°C) as the working fluid in the differential manometer. The nitrogen diluent flowrate was measured using a Poiseuille flowmeter described previously.⁸

The diluent nitrogen was either prepurified nitrogen obtained from the Matheson Company or white spot nitrogen from the Air Products Company. Technical grade (98.5%) NO, $\text{Ni}(\text{CO})_4$, and C_2H_2 (99.6%) were obtained from the Matheson Company. CO (99.3%) was obtained from the Air Products Company. NO was purified by vacuum distillation from a trap maintained at 113°K to one cooled with liquid nitrogen, with initial and final fractions being discarded. The $\text{Ni}(\text{CO})_4$ was purified by freezing down in a liquid nitrogen trap and pumping with a mechanical pump and then a Hg diffusion pump. The liquid nitrogen was then removed and the $\text{Ni}(\text{CO})_4$ was distilled to a trap at -77°C

with initial and final fractions being discarded. The Ni(CO)_4 was a colorless liquid at room temperature. C_2H_2 was purified by vacuum distillation from a trap maintained at 196°K to one cooled with liquid nitrogen. CO was passed through a glass wool-packed U-trap at 78°K and was added directly to a storage bulb. Gases were always conducted from the tanks in glass tubing. Glass-metal connections were all made either with soldered Kovar seals or with epoxy adhesive.

As the Ni(CO)_4 -active nitrogen reaction proceeded a black deposit occurred on the downstream walls of the flowtube and on the mouth of the Ni(CO)_4 inlet. Coating the walls of the flowtube with 85% phosphoric acid saturated with phosphorus pentoxide allowed at least 5 hours of operation before buildup of the deposit on the walls became objectionable.

All measurements were taken at room temperature, which averaged $22 \pm 3^\circ\text{C}$.

RESULTS

Description of the Flame

The addition of Ni(CO)_4 to a flow of discharged nitrogen produced a conical, bright blue flame which extended down the flow tube. The brightest flames were obtained for runs around 3 torr. For $[\text{Ni(CO)}_4]$ on the order of $[\text{N}]$ the flame was typically observed for at least 15 cm downstream. Under these conditions the yellow afterflow was also visible downstream of the carbonyl inlet. Increasing the flow of the Ni(CO)_4 increased the width of the conical flame, decreased the extent to which it extended downstream, and decreased the intensity of the yellow afterglow downstream. As the pressure was decreased to 0.3 and 0.13 torr the blue flame became very diffuse. At 0.14 torr the flame extended about 3 cm upstream and 10-15 cm downstream. At these low pressures

it was necessary to increase the length of the sidearm to which the fused silica window was attached. Otherwise, the flame extended out to the window and a black deposit eventually coated the window so extensively the observed intensity was drastically reduced although the flame was approximately constant. As the pressure in the flow tube was increased to 9.6, 14.7, and 27 torr the flame began to develop what appeared to be a reddish cone inside the blue cone.

On occasion, especially with added O_2 , when the $Ni(CO)_4$ was shut off and only the yellow nitrogen afterglow was present, a distinct red glow was present on the walls of the flow tube where the black deposit was present. This was assumed to be the surface catalyzed excitation of nitrogen that has been studied by Harteck, Mannella, and Reeves.⁹ Since there is also a (nickel) deposit at the carbonyl inlet, it is possible that the reddish inner cone observed at higher pressures may be due to this phenomenon.

Pressure Dependence of the Atomic Nickel Emission Spectra

The atomic nitrogen-nickel carbonyl flame consists almost entirely of emission from excited nickel atoms. The only other emission observed was weak CN violet emission. No other emission has been detected in the wavelength range from 2300 to 7600Å. Table I lists most of the stronger lines, the energy level of the emitting state and its term symbol, and the relative intensities of the lines at a representative number of pressures between 0.13 and 27 torr. The type of discharge, pulsed (P) or microwave (M), is indicated in the parentheses next to the pressure. The data for 0.13 and 0.54 torr had 0.6% O_2 added before the discharge to promote dissociation on the nitrogen. All intensities are recorded relative to an arbitrarily chosen standard peak at 3369.57Å which has been assigned an intensity of 500. Over the several hours necessary for one run the flame intensity was often found to decrease

by about 15-30% due to clogging of the carbonyl inlet with a solid deposit. In all runs the standard peak was swept repeatedly, any changes in flame intensity were assumed to be linear, and the observed relative intensities were adjusted to standard conditions by linear interpolation. Results at a given pressure are considered reproducible to within 10-15%.

Experiments were run using a microwave discharge for runs greater than 1 torr and a pulsed discharge for runs less than 1 torr. Data taken with both the pulsed discharge and the microwave discharge at 1.0 torr gave identical results within scatter.

Wavelength assignments were made by comparing the photoelectric spectra with the compilation of atomic nickel lines by Corliss.¹⁰ Line identifications were verified between 2300 and 5200 Å by photographic spectra taken with 50 μ slits (0.5 Å bandpass) and analyzed on a measuring microscope.

Reactant Ratio

The nitrogen atom and $\text{Ni}(\text{CO})_4$ concentrations used in the reported data are summarized in Table II. One experiment was performed to determine if the reactant ratio affected the observed relative intensities. At 1.8 torr five runs were taken in which the reactant ratio, $[\text{N}]/[\text{Ni}(\text{CO})_4]$, was varied between 4.5 and 29 and representative strong lines over the entire wavelength range were recorded. Within normal fluctuation the relative intensities were not affected by changing reactant ratio.

Observation Time

The linear flowrate of the diluent nitrogen varied for the different pressure runs between about 100 cm/sec and 550 cm/sec. The higher pressure runs had lower flowrates due to throttling and the lower pressure runs had lower flowrates due to decreased pumping efficiency. Since the carbonyl inlet was stationary, slightly different portions of the atomic flame were observed.

Two experiments were performed to determine if the observation time after mixing had any affect on the observed relative intensities. At 1.8 torr two runs were taken with linear flowrates of 100 and 550 cm/sec. The relative intensities were the same for the two runs. For all runs mentioned thus far the Ni(CO)_4 was added through the inlet at the observation window. Two runs were performed in which Ni(CO)_4 was added through the upstream inlets. With the Ni(CO)_4 added through the inlet 19 cm upstream at a total pressure of 3.6 torr the observed intensities were very weak, and thus not well determined, but within the accuracy of the data the relative intensities were not affected by the inlet position. At 9.5 torr Ni(CO)_4 was added through the inlet 5.3 cm upstream (equivalent to 4×10^{-2} sec at the prevailing flowrate) and compared to the run with Ni(CO)_4 added at the observation window. Within scatter there was no change in the relative intensities. The observed relative intensities are not affected by changes in observation time to the extent which they are varied in these experiments.

Effect of Added Gases on Relative Intensities

a) Oxygen added before the discharge

Oxygen when added to the nitrogen flow before the pulsed discharge has no appreciable affect on the relative intensities. The addition of 0.6% O_2 to a run with the pulsed discharge at 0.3 torr gave relative intensities about 20% lower than a run without added O_2 . Thus the 0.13 torr run with 0.6% O_2 recorded in Table I is not significantly affected by the added O_2 .

With O_2 added before the microwave discharge there was an effect, though small. A run at 1.85 torr with 0.3% O_2 had all peaks increased by an average of about 20% relative to the standard peak. For a run at 1.1 torr with 0.6% added O_2 lines with $E_u > 40,000 \text{ cm}^{-1}$ increased about 15%. A comparison for lower pressures can be made using the data in Table I. It was

not possible to perform low pressure experiments with a microwave discharge and no added O_2 because of the forbiddingly low nitrogen atom concentration and consequent weak atomic flame. However, the pulsed discharge experiment at 0.3 torr can be compared with the microwave discharge experiments at 0.4 torr with 0.6% O_2 . In both cases two runs giving very similar results were averaged to give the reported data so the accuracy is fairly good. Again the microwave discharge with added O_2 gives higher relative intensities for lines with $E_u > 40,000 \text{ cm}^{-1}$. For lines with $E_u < 40,000 \text{ cm}^{-1}$ the increase in relative intensity with added O_2 is only slightly greater than normal scatter. However, lines originating from the higher energy states had increases in relative intensity that were larger than reasonable scatter. It does appear that these higher energy states are increased in intensity relative to the lower states. A possible explanation is that the lower states or their precursor are quenched by O_2 or a discharge product of O_2 . Since the effect is not observed with a pulsed discharge it would appear that a microwave discharge product of O_2 is involved. Considering, however, the marginal extent of the observed effect, very little can be said in this matter.

b) Added C_2H_2

For a run at 1.8 torr with a partial pressure of added $Ni(CO)_4$ of 0.9×10^{-3} torr, C_2H_2 was added in four differing amounts between 0.9×10^{-3} torr and 1.2×10^{-2} torr. There was no effect on the relative intensities of a representative sampling of lines. For a run at 5.6 torr 4×10^{-3} torr of C_2H_2 had no effect on a flame with 1.7×10^{-3} torr added $Ni(CO)_4$.

c) Added CO

The addition of CO to the carbonyl flame at 2.1 torr in concentrations between 3 to 20 times that of $Ni(CO)_4$ had no effect on the observed relative intensities.

Trap Analysis

The black deposits taken from the carbonyl inlet and from the U-traps downstream of the reactor were analyzed for nickel content using either a titration with sodium ethylenediaminetetraacetate (EDTA)¹¹ and murexide indicator or precipitation with dimethylglyoxime (DMG)¹¹. The results are summarized in Table III. The deposit in the inlet is almost entirely nickel metal. The deposit collected in the U-traps appears to vary in nickel content between 35-70%. The trap deposit may consist of nickel plus a polymeric compound or nickel as part of a polymeric compound. Frequently, sparks were observed in the U-traps when air was introduced, so apparently at least a portion of the deposit is some form of active nickel "black", which is pyrophoric.

CN Violet Band System

CN violet ($B^2_{\Sigma^+} \rightarrow X^2_{\Sigma^+}$) emission has been observed weakly in the carbonyl flame at low pressures and with a moderate intensity in the high pressure flames. Some of the observed bands at several pressures are summarized in Table IV. The CN emission is present only when $Ni(CO)_4$ is added; it is not an impurity emission.

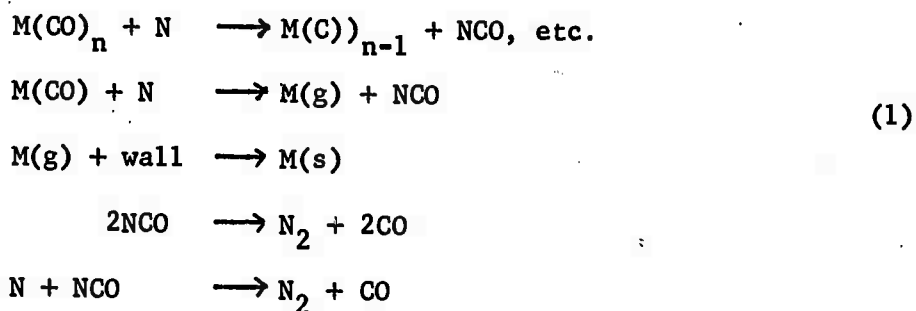
The study of CN emission in flames is a very complex matter with several different proposed mechanisms of CN excitation^{12-17,4}. A reaction mechanism proposed for the active nitrogen-nickel carbonyl reaction by Brennen and Kistiakowsky⁵ included CO as a reaction product and NCO as an intermediate. The reaction of atomic nitrogen with CO to form ground state CN is 74 kcal/mole endothermic and will not be important. The reaction, $N + NCO \rightarrow CN + NO$, is 7 kcal/mole exothermic and may provide a source of ground state CN. The CN could then be excited to its B^2_{Σ} state by collision with $N_2(A^3_{\Sigma_u^+})$, which is present in the active nitrogen. Such a mechanism has been proposed by

Bayes¹² and supported by Campbell and Thrush¹⁵. The similar CN vibrational level distribution found at 9.6, 14-7, and 27 torr suggests that just one mechanism is operating at these pressures.

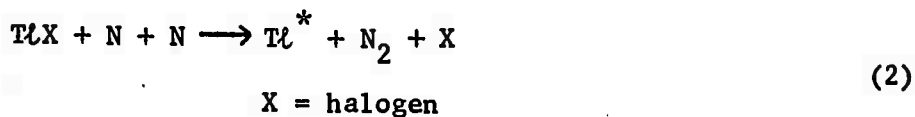
DISCUSSION

Excitation of Metal Atoms by Active Nitrogen

Since the earliest observations of Strutt¹ that the line spectra of various metals - cadmium, magnesium, mercury, potassium, zinc, and lead - are developed when the metals are exposed to active nitrogen, several studies²⁻⁵ of excitation of metal atoms in active nitrogen have been made. Brennen and Kistiakowsky⁵ (BK) found that $\text{Ni}(\text{CO})_4$, $\text{Fe}(\text{CO})_5$, $\text{Cr}(\text{CO})_6$, $\text{W}(\text{CO})_6$, $\text{Mn}_2(\text{CO})_{10}$, and $\text{Co}(\text{NO})(\text{CO})_3$ reacted rapidly with active nitrogen, forming metal atoms, and producing flames due to metal atom emission spectra. Studying the $\text{Ni}(\text{CO})_4$ and $\text{Fe}(\text{CO})_5$ reactions in detail BK concluded that a stepwise degradation occurred

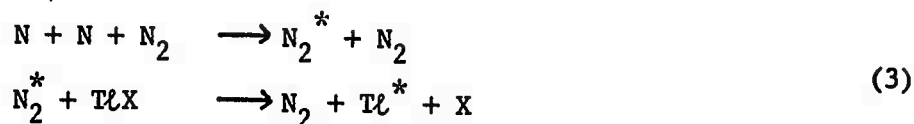


where M is the metal atom. The emission spectra was concluded to result from excitation of metal atoms in collision with metastable $\text{N}_2(\text{A}^3\Sigma_u^+)$ molecules. Phillips^{2,3}, observing primarily the thallous halides, concluded that the thallium atom was excited in a termolecular process.



or in a bimolecular reaction of TlX with an excited nitrogen molecule formed

during N-atom recombination



Since the Tl atoms were assumed to be formed as the result of a dissociative collision process, it was necessary for N_2^* to have at least 225 kcal/mole of energy to form Tl in its observed states. $\text{N}_2(^5\Sigma)$ was considered to be the most probable excited nitrogen molecule. It is possible that the reaction of the thallous halides with active nitrogen also involves stripping of the halide atom in the reaction of TlX with a N-atom and consequent excitation of the free Tl-atom in collision with a $\text{N}_2(\text{A})$ molecule. The observations led Phillips to discount a degradation mechanism: (1) no Tl emission was found in the region immediately following the flame, and (2) the addition of Tl vapor alone to active nitrogen resulted in a much weaker flame. If in either of these cases the Tl-atom concentration were significantly lower than in the TlX flame, a much weaker emission flame would be expected. The addition of Tl vapor, indeed, suggests that the Tl-atom can be directly excited by active nitrogen.

Our observations of the $\text{Ni}(\text{CO})_4$ flame are similar to those of BK^5 . We found that the relative intensities of the various nickel lines is the same at the reaction inlet and at positions much further downstream. This supports the proposed degradation mechanism of BK with excitation of gaseous nickel atoms by $\text{N}_2(\text{A})$ which is present in the active nitrogen downstream of the reaction flame. We have observed excited nickel atoms with up to about $57,000 \text{ cm}^{-1}$ of energy above the ground state, in agreement with BK. Our observations were limited primarily to states with $E_u < 52,000 \text{ cm}^{-1}$. If Ni is assumed to be excited from its ground state, then $\text{N}_2(\text{A})$ molecules in vibrational levels up to at least $v' = 6$ are necessary to supply the required

excitation energy. Ni has low lying states which may be populated to some degree. An absorption study of the carbonyl flame is necessary to determine if the low lying states are appreciably populated. If they are, then slightly lower vibrational levels of $N_2(A)$ could supply the necessary excitation energy.

Pressure Dependence of the Relative Steady-State Population Distribution

Relative steady-state populations for the excited nickel levels were calculated from the relative intensities recorded in Table I. The relative intensities were corrected for the relative sensitivity of the photomultiplier tube and monochromator, Fig. (3), and populations were calculated from the relationship

$$N_u \propto \frac{g_u I \lambda^2}{g_l f_{lu}} \quad (4)$$

where N_u is the steady-state population, I is the relative intensity in units of quanta, λ is the wavelength, f_{lu} is the absorption oscillator strength, and g_u and g_l are the degeneracies of the upper and lower states of a transition, respectively. g -Values were taken from Corliss¹⁰. Plots of relative population against energy of the upper state are given in Fig. (4) through Fig (11) for pressures between 0.13 and 27 torr. A line representing an effective Boltzmann temperature is drawn through the data. Considering the scatter in the data the temperature is accurate to about $\pm 400^\circ K$. A significant change in the relative population distribution is found over the pressure range studied as demonstrated by the change in effective temperature. A plot of the observed effective temperature against pressure is given in Fig. (12).

For purposes of qualitative evaluation of trends the nickel population has been divided into four groups with different E_u : (1) group I, $28,000 < E_u < 31,000 \text{ cm}^{-1}$; (2) group II, $31,000 < E_u < 37,000 \text{ cm}^{-1}$; (3) group III, $42,000 < E_u < 45,000 \text{ cm}^{-1}$; and (4) group IV, $45,000 < E_u < 52,000 \text{ cm}^{-1}$. We have taken only limited data for energy levels above $52,000 \text{ cm}^{-1}$. Observation

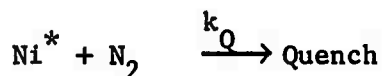
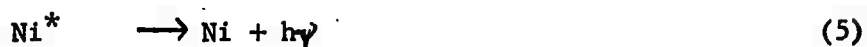
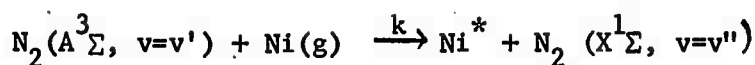
of Fig. (4) through Fig. (11) reveals that in addition to the overall change in relative population there has been a noticeable change in the relative position of group II. The effective temperature line has been drawn largely through groups I, III, and IV. At lower pressures the group II points fall below this temperature line while at the higher pressures they have increased in population and fall very nicely in line. Between 0.13 and 27 torr group II has risen from an average of about 25% lower in population than group I to a population very nearly equal to group I.

With two exceptions the relative populations within a given group did not change with pressure. That is, the higher population points within group IV at 0.13 torr and 27 torr represent the same emission lines, etc. The 3371.99A and 2994.46A lines violate this trend and will be discussed below in the section on quenching.

The lower energy levels of Ni can be populated by radiative cascading from higher levels. Relative intensities for data at 9.6 torr were converted to intensities proportional to quanta/sec and the ratio of cascading from higher levels, $E_u = 42-52,000 \text{ cm}^{-1}$, to radiative depletion was calculated. For levels with $26,000 < E_u < 32,000 \text{ cm}^{-1}$ the ratio was 0.13, and for levels with $32,000 < E_u < 37,000 \text{ cm}^{-1}$ the ratio was 0.16. An estimate was also made including lines listed in "Tables of Spectral-line Intensity" by Meggers, et al.¹⁸ that we did not observe directly. The ratios were only minimally changed by this correction. It can be seen from the low ratios that the population of the lower levels is not affected appreciably by radiative cascading.

Excitation Mechanism

A mechanism which describes the excitation of the Ni-atoms is as follows:



Excited nickel atoms, Ni^* , are formed in collision with the metastable $\text{N}_2(\text{A})$ molecule which drops to its ground state. The excited nickel atoms are depopulated by radiation or quenching by nitrogen buffer. The population distribution in Ni^* reflects the energy distribution in the vibrational levels of $\text{N}_2(\text{A})$ and the rate of population of the different excited states of nickel. If the quenching step is neglected (discussion below of the quenching step suggests this is a reasonable assumption) the rate of formation of a given energy level, E_u , is represented as

$$R = N_u / \tau_u \quad (6)$$

where N_u is the steady-state population and τ_u is the radiative lifetime. Using relative populations and calculated radiative lifetimes (see below, Table VI), relative rates of formation of some of the states with strong emission lines have been evaluated at 0.13, 2.4, and 27 torr. See Figs. (13) to (15).

An attempt was made to calculate an expected rate of formation of Ni^* as a function of E_u using a known vibrational distribution for $\text{N}_2(\text{A})$. An important issue in this calculation is the importance of the Franck-Condon principle for the $\text{N}_2(\text{A}) \rightarrow \text{N}_2(\text{X})$ transition in determining the energy that is available in a bimolecular collision with $\text{Ni}(\text{g})$. Robertson¹⁹ in a study of collision of excited Ar atoms with molecules has concluded that the molecules do obey the Franck-Condon principle. Transitions in which molecular separations are maintained constant are the most favored. Callear and Wood²⁰ have considered the energy transfer reactions of $\text{N}_2(\text{A})$ with a series of other molecules. They concluded that to the first approximation of $\text{N}_2(\text{A}) \rightarrow \text{N}_2(\text{X})$ transition obeyed the Franck-Condon principle.

In a separate study²¹ the $N_2(A)$ vibrational population distribution was determined under a variety of conditions. Table V is a summary of the observed relative vibrational populations at 2.4 and 40 torr. The relative rates of formation were calculated according to

$$\frac{d(Ni^* \text{ (group I)})}{dt} \propto k(q_{v',v''}) [N_2(A, v=v')][Ni(g)] \quad (7)$$

for all transitions in the range $E(A, v=v') - E(X, v=v'') = E \text{ (group I)}$, etc. The rate of energy transfer from $N_2(A)$ to $Ni(g)$ was considered to be of the form¹⁹

$$k = v [\sigma P_E(\Delta E)] q_{v',v''} \quad (8)$$

where v is the relative velocity of the collision partners, σ is the reaction cross section, $P_E(\Delta E)$ is the potential energy part of the cross section, and $q_{v',v''}$ is the Franck-Condon factor. For calculation of relative rates of formation the cross sections are all assumed to be equal so that

$$k \propto q_{v',v''} \quad (9)$$

Several $N_2(A)$ vibrational levels feed into a given energy range so the relative rates of formation were summed according to group. $[q_{v',v''}] [A_{v'}]$ was evaluated for various v' 's and summed for all transitions with the energy for the $N_2(A) \rightarrow N_2(X)$ transition within the energy range for group I, etc. The F-C factors were taken from Benesch et al.²². The calculated relative rates of formation are shown as squares in Fig. (14) and Fig. (15). The relative rates are recorded at an average energy for a given group. There is good agreement between the calculated and observed relative rates of formation. The agreement indicates that to a first approximation the proposed mechanism adequately describes the excitation of nickel atoms. The assumption that the Franck-Condon principle applies for the $N_2(A) \rightarrow N_2(X)$ transition in collision

appears to be valid. Indeed, the nickel carbonyl-active nitrogen reaction provides a good test of the applicability of the Franck-Condon principle in bimolecular collisions. In the energy transfer reactions of $N_2(A)$ with other molecules²¹ only a limited number of acceptor states are available. With the Ni-atom a whole manifold of acceptor states are available. It is desirable to look at the excitation of other metals, with reasonably extensive energy manifolds, by $N_2(A)$ to determine if the Franck-Condon principle is generally applicable in collisions. The reaction of other metal carbonyls with active nitrogen provides a very convenient method for studying this process.

The change in the relative population of the group II lines is a detail we have not clarified. The calculated rates of formation are not sufficiently accurate to determine if the change reflects the change in the $N_2(A)$ vibrational distribution or if another factor is involved. Figs. (13) to (15) suggest that group II does increase relative to group I as a result of the change in the vibrational distribution of $N_2(A)$.

The Quenching Step

In order to evaluate the importance of the quenching step, equation (5), relative intensities of long-lived states, the states expected to be most readily affected by the quenching reaction, were analyzed at different pressures. The radiative lifetimes of a number of the excited states were calculated according to the relationship

$$\tau^{-1} = g_l^{-1} \sum_l g_l A_{ul} \quad (10)$$

and are summarized in Table VI. The g-values and gA values were taken from Moore²³ and Corliss and Bozman²⁴, respectively. Steady-state analysis of the mechanism, equation (5), gives

$$I_{\text{rel}}(Ni^*)^{-1} = \tau^{-1}/R + k_Q[N_2]/R \quad (11)$$

where R is the rate of formation of Ni^* , and $I_{\text{rel}}(\text{Ni}^*) \propto [\text{Ni}^*]$. A plot of $I_{\text{rel}}(\text{Ni}^*)^{-1}$ against pressure will have a slope to intercept ratio of $k_Q \tau$. k_Q can be evaluated using the calculated values of τ .

The 3371.99A line originating from the long-lived 30980 cm^{-1} state is the one observed line which is drastically affected by pressure. This line, which is labeled in Fig. (4) through Fig. (11), drops in relative population by a significant degree between 0.13 and 27 torr. A graph of the reciprocal relative intensity, $I(3371.99)^{-1}$, against pressure is shown in Fig. (16). A value of $k_Q = 3 \times 10^{-10}$ is calculated using $\tau = 2 \times 10^{-8}$ sec. The data is very scattered but the quenching mechanism appears at least approximately correct. The calculated quenching coefficient is quite high, being essentially gas kinetic, i.e., the quenching reaction occurs on every gas kinetic collision. None of the lines originating from the other long-lived states showed any significant change. Lines originating from the 29833 cm^{-1} state, $\tau = 6 \times 10^{-9}$ sec, did not exhibit a change in relative population. The 2994.46A line, see Figs. (4) to (11), which originates from the 33590 cm^{-1} state, $\tau = 5 \times 10^{-9}$ sec, did exhibit a slight drop in relative population. It can be seen from Table VI that both the 30980 cm^{-1} and 33590 cm^{-1} states, in addition to having longer lifetimes, have nearby lower energy states within 100 cm^{-1} while the next lowest energy level below the 29833 cm^{-1} state is 400 cm^{-1} away. (At room temperature $kT \cong 200 \text{ cm}^{-1}$). It appears, meager though the evidence is, that both a long radiative lifetime and a nearby acceptor state are necessary for the quenching step to be important. Ni^* is quenched to its next lower electronic state and N_2 absorbs the energy as an increase in translational energy. States with radiative lifetimes shorter than 5×10^{-9} sec were not affected by the quenching step.

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TABLE I: SUMMARY OF RELATIVE INTENSITIES

λ (Å)	E_u (cm ⁻¹)	Term Symbol	Relative Intensity							
			(P) 0.13 torr.	(P) 0.31 torr.	(M) 0.4 torr.	(M) 1.1 torr.	(M) 2.3 torr.	(M) 9.6 torr.	(M) 14.7 torr.	(M) 27 torr.
2313.98	45419	³ wF ₂	-	-	60	-	118	91	(44)	70
2325.79	44315	³ yG ₄	-	-	63	-	-	-	-	-
2337.82	44475	³ wD ₂	14	-	86	-	113	72	40	42
2337.49	42728	³ wD ₃								
2424.03	42954	³ xD ₂	19	-	72	-	73	34	16	18
53.99	42954	³ xD ₂	(8)	-	30	-	-	10	-	-
72.06	42656	³ yP ₁	18	-	55	-	31	25	-	-
76.87	40361	¹ o ₃	52	-	130	-	45	25	-	-
2981.65	34409	³ yD ₁	105	115	90	145	85	250	330	520
94.46	33590	¹ zG ₄	1050	1100	950	920	700	570	560	580
3002.48	33501	³ yD ₃	1550	1950	1700	2100	1900	3250	3900	5000
3.63	34163	³ yD ₂								
12.00	36601	¹ yD ₂	660	950	740	1000	1100	1600	2300	3100
19.14	33112	³ yF ₃	51	-	54	85	-	-	110	(56)
37.94	33112	³ yF ₃	540	500	520	(650)	680	910	1000	1200
50.82	32973	³ yF ₄	2250	2700	(1950)	2900	2600	2600	2700	2700

$\lambda(\text{\AA})$	$E_u(\text{cm}^{-1})$	Term Symbol
54.31	33611	3^3F_2
57.64	34409	y^3D_1
64.62	33501	y^3D_3
80.76	34163	y^3D_2
99.12	33590	z^1G_4
3101.56	33112	y^3F_3
3101.88	35639	y^1F_3
14.12	32982	z^1P_1
34.10	33611	y^3F_2
3213.41	58526	e^5P_1
3214.06	56858	e^5D_4
3215.82	57744	e^5D_3
3217.32	56821	e^5P_3
3284.42	57104	i^3D_3
3369.57	29669	z^3D_3
3371.99	30980	z^3G_4
3374.22	29833	z^5F_3
74.64	56885	e^5H_7
3380.57	32982	z^1P_1
91.04	29481	z^3F_4
92.98	29669	z^3D_3
3414.76	29481	z^3F_4
23.71	30913	z^3D_1

Relative Intensity									
(P)	(P)	(M)	(M)	(M)	(M)	(M)	(M)	(M)	(M)
0.13 torr.	0.31 torr.	0.4 torr.	1.1 torr.	2.3 torr.	9.6 torr.	14.7 torr.	27 torr.		
480	420	360	670	710	900	1000	1300		
-	520	470	770	780	1300	1700	2600		
130	140	130	195	250	290	390	450		
100	115	90	150	85	130	200	210		
(125)	(240)	190	210	113	-	-	-		
2050	2500	2100	2900	2600	3100	3300	3600		
(40)	43	25	48	-	65	(70)	110		
620	680	680	-	1100	1500	1900	1900		
9	-	-	-	-	-	-	-		
16	-	-	-	-	-	-	-		
40	49	-	-	-	-	-	-		
500	500	500	500	500	500	500	500		
510	420	440	200	160	72 \pm 15	65 \pm 15	50 \pm 15		
-	-	(90)	(85)	59	-	(16)	-		
680	630	590	560	670	690	790	1050		
-	-	(130)	-	220	-	-	-		
650	700	660	680	690	730	690	700		
1750	2000	1600	1700	1650	2100	1700	1800		
340	320	340	350	380	280	250	260		

$\lambda(\text{\AA})$	$E_u(\text{cm}^{-1})$	Term Symbol
3433.56	29321	3^3F_3
37.28	29084	5^5F_4
46.26	29888	3^3D_2
52.89	29833	5^5F_3
58.47	30619	3^3F_2
61.65	29084	5^5F_4
67.50	30163	5^5F_2
72.54	30913	3^3D_1
83.87		
92.96	29501	3^3P_1
3515.05	29321	3^3F_3
24.54	28569	3^3P_2
66.37	31442	1^1D_2
71.87	29321	3^3F_3
87.93	28068	5^5G_4
3619.39	31031	3^1F_3
3783.53	29833	3^5F_3
4401.54	48467	5^5F_5
35.35	52041	3^3F_2
4459.03	49086	5^5F_4
62.45	50346	5^5F_2
66.38	52272	3^3D_2

Relative Intensity										
(P)	(P)	(M)	(M)	(M)	(M)	(M)	(M)	(M)	(M)	(M)
0.13 torr.	0.31 torr.	0.4 torr.	1.1 torr.	2.3 torr.	9.6 torr.	14.7 torr.	27 torr.			
320	290	400	330	370	310	260	240			
200	270	210	(170)	220	93	100	110			
760	690	850	720	930	860	760	730			
540	440	510	390	460	280	270	260			
620	800	750	610	770	(650)	640	590			
1450	1400	1500	1000	1400	(850)	800	720			
85	-	(70)	-	-	-	-	-			
320	300	330	280	400	280	330	370			
120	140	140	110	113	85	100	110			
600	530	700	500	900	(700)	690	730			
1100	800	1100	770	1350	850	680	(600)			
1650	1500	2000	1250	2300	1850	1400	1400			
680	580	880	(530)	1200	1200	1100	1250			
100	110	130	80	110	85	-	-			
180	100	120	(38)	-	-	-	-			
940	860	1200	800	1600	1200	860	880			
75	-	120	-	85	60	46	52			
330	350	430	270	380	300	200	150			
38	-	58	29	25	(10)	-	-			
190	190	260	150	208	110	58	47			
-	-	75	44	51	42	(28)	-			
-	-	40	(25)	6	-	-	-			

$\lambda(\text{\AA})$	$E_u(\text{cm}^{-1})$	Term Symbol	Relative Intensity									
			(P)	(F)	(M)	(M)	(M)	(M)	(M)	(M)	(M)	(M)
			0.13 torr	0.31 torr	0.4 torr	1.1 torr	2.3 torr	9.6 torr	14.7 torr	27 torr		
4470.43	49778	$5F_3$	95	95	140	84	100	73	34	-		
4600.36	50745	$5F_1$	32	-	70	37	-	24	-	-		
5.00	49778	$5F_3$	74	100	-	85	125	64	44	35		
48.65	49086	$5F_4$	135	185	300	150	85	100	85	49		
86.21	50346	$5F_2$	(30)	-	-	34	26	(29)	22	23		
4703.81	50754	f^1D_2	50	-	(80)	40	37	29	(22)	-		
14.42	48467	$5F_5$	490	430	700	340	490	350	290	250		
52.12	50706	$1G_4$	47	-	85	35	46	10	(13)	-		
52.42	50537	$1P_1$										
56.51	40986	$5F_4$	100	140	200	93	145	64	47	35		
63.93	50466	$3F_4$	40	-	90	39	-	21	(15)	(11)		
86.54	48467	$5F_5$	145	180	220	115	171	125	105	83		
4808.89	50678	$3G_3$	(58)	-	90	-	12	32	(19)	18		
29.02	49271	$3D_3$	(160)	150	(260)	135	171	83	53	-		
38.64	54251	$3F_3$	-	-	45	25	32	24	-	-		
55.41	49159	$3G_5$	350	290	490	310	270	120	80	53		
66.27	49086	$5F_4$	125	160	220	100	151	90	46	43		
73.44	50346	$5F_2$	67	(90)	110	57	76	55	34	34		
4904.41	48953	$3S_1$	280	250	430	180	240	120	65	46		
13.97	51306	$3F_3$	54	-	80	40	47	32	(17)	26		

$\lambda(\text{\AA})$ $E_u(\text{cm}^{-1})$ Term Symbol

Relative Intensity

		(P)	(P)	(M)	(M)	(M)	(M)	(M)	(M)	(M)
		0.13 torr	0.31 torr	0.4 torr	1.1 torr	2.3 torr	9.6 torr	14.7 torr	(M)	(M)
4918.36 } 4918.71 } 4925.56 } 4935.83 } 37.35 } 53.21 } 65.17 } 71.34 } 80.17 } 84.11 }	51306	100	155	230	110	148	95	63	46	27
	50717	25	-	(30)	23	29	18	(10)	19	
	51306	130	110	190	(93)	96	70	39	39	
	52040	45	(60)	90	37	35	35	(28)	27	
	49333	12	-	20	-	-	10	-	-	
	50346	62	80	125	60	62	49	(31)	32	
	50754	400	350	550	240	320	135	93	61	
	56711	320	310	500	230	300	115	81	43	
	49158									
	50678									
5003.74 } 17.58 } 18.29 } 35.36 } 35.96 }	33501	135	170	230	100	108	48	(41)	26	
	48467	320	290	440	200	310	200	170	130	
	50834	800	690	1100	530	690	250	180	130	
	49175									
	49333									
48.84 } 80.53 } 81.11 }	50832	(100)	160	190	84	90	50	32	-	
	49158	1350	1400	2000	880	1000	470	320	220	
	50706									
84.09 } 99.31 } 99.93 }	49333	470	(380)	(600)	320	410	-	(110)	86	
	49086	450	330	(500)	280	345	170	110	72	
	49271									

λ (Å)	E_u (cm ⁻¹)	Term Symbol
5115.39	50466	³ F ₄
25.23	49175	³ G ₄
29.37	49159	³ G ₅
42.78	49328	³ D ₃
46.48	49314	³ F ₃
55.12	50834	³ F ₂
55.76	50832	¹ F ₃
68.66	49175	³ G ₄
6176.81	49158	³ G ₅
7122.20	42606	³ D ₃
7393.60	42606	³ D ₃
7409.35	44112	³ D ₁
22.28	42790	³ D ₂
7555.60	44263	¹ D ₂
74.05	44112	³ D ₁

Relative Intensity

(P)	(P)	(M)	(M)	(M)	(M)	(M)	(M)	(M)
0.13 torr	0.31 torr	0.4 torr	1.1 torr	2.3 torr	9.6 torr	14.7 torr	27 torr	(M)
280	260	(350)	200	250	120	87	48	
(60)	-	110	56	60	27	(24)	(11)	
90	(125)	160	77	66	44	(27)	28	
280	210	320	170	160	-	63	53	
320	340	500	240	290	140	78	68	
240	220	340	180	180	-	53	76	
115	125	180	90	90	-	25	25	
-	-	-	-	60	41	-	-	
135	170	250	100	118	51	47	47	
21	-	75	38	-	20	16	11	
53	-	60	41	-	24	19	19	
48	-	90	55	-	26	24	20	
-	-	85	-	-	32	24	21	
-	-	(20)	-	-	12	(8)	-	

Table II:

Summary of Reaction Conditions

Pressure	Discharge ^a	[N]	[Ni(CO) ₄]	[N]
				[Ni(CO) ₄]
		x10 ³ , torr	x10 ³ , torr	
0.13	P	3.5	0.5	7.2
0.31	P	1.3	0.6	2.1
1.1	M	0.7	0.18	3.9
2.3	M	7.6	0.56	14
9.6	M	5.1	1.1	4.6
14.7	M	7.0	0.47	15
27	M	13	3.6	3.6

^aP represents a pulsed discharge and M a Microwave discharge.

Table III:
Analysis of Solid Deposits

Source	Method ^a	%Ni
U-trap	EDTA	36
U-trap	EDTA	35
U-trap	DMG	70
inlet	EDTA	97
inlet	EDTA	95
inlet	DMG	92
inlet	DMG	94

^a Titration with sodium ethylenediaminetetraacetate (EDTA) or precipitation with dimethylglyoxime (DMG).

Table IV:
CN Violet ($B^2_{\Sigma} \rightarrow X^2_{\Sigma}$) Relative Intensity

CN Band (v', v'')	Relative Intensity ^a				
	0.3 torr	0.4 torr	9.6 torr	14.7 torr	27 torr
3,2					320
2,1					270
1,0					210
4,4	(11)	(14)	210	140	490
3,3			155	95	320
1,1			137	85	310
0,0	32	21	210	140	510
5,7				28	95
0,1				28	42

^aInaccurately determined values are in parentheses.

Table V:

Relative Vibrational Population of $N_2(A^3\Sigma^+)$

$N_2(A^3\Sigma^+)$ v'	Relative Population	
	2.4torr	40torr
0	130	200
1	130	200
2	120	50
3	42	16
4	18	5

Table VI:

Calculated Radiative Lifetimes for Excited States of Nickel

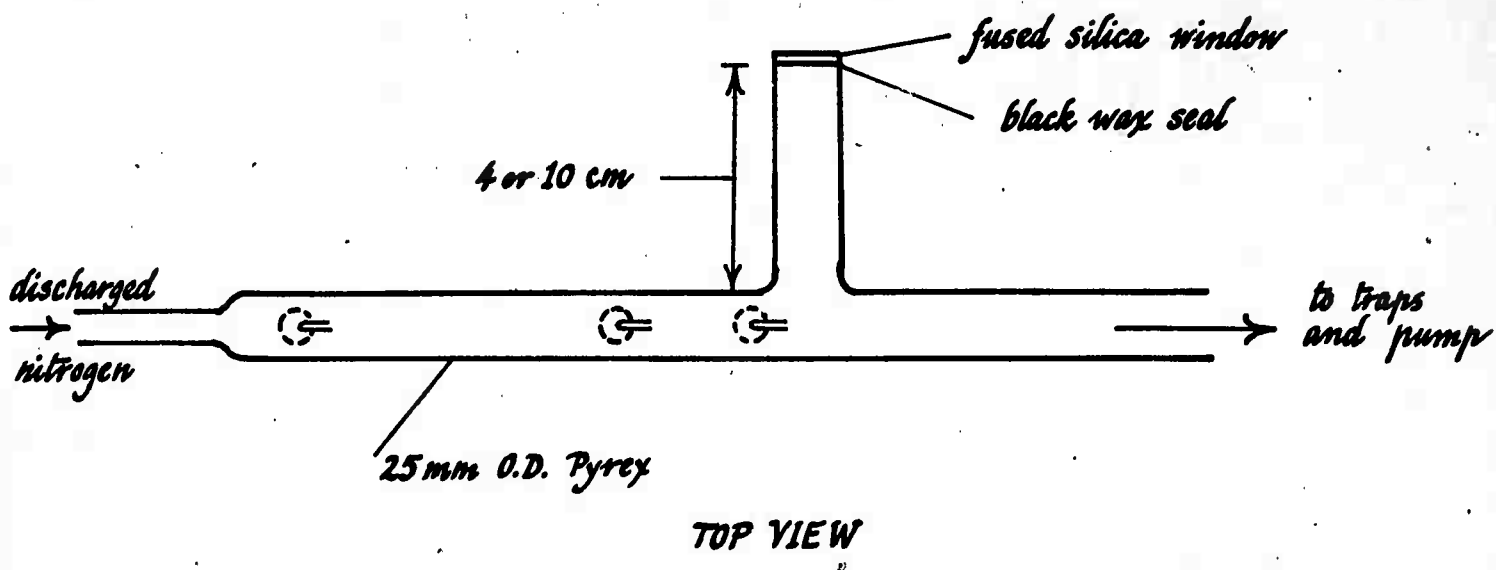
E_u, cm^{-1}	$\tau \times 10^9, \text{sec}^a$	E_u, cm^{-1}	$\tau \times 10^9, \text{sec}^a$
28569	2.7	42790	1.3
29084	3.4	44112	(3.8)
29321	1.3	48467	0.34
29481	2.2	48953	2.3
29669	2.0	49086	0.23
29833	6.4	49158	0.37
29888	1.9	49159	0.88
30913	2.5	49175	0.41
30980	21	49271	1.3
31031	1.1	49314	(0.89)
31442	1.5	49328	(2.0)
32973	2.3	49333	(1.9)
32982	1.5	49778	0.37
33112	1.5	50346	0.33
33501	1.3	50466	0.70
33590	4.9	50717	0.51
33611	0.79	50832	(0.66)
34163	1.7	51306	1.3
34409	0.8	52040	(2.0)
35639	1.8		
36601	0.66		
42606	1.4		

^aCalculated from $\tau^{-1} = g_u^{-1} \sum_l g_l A_{ul}$. For values in parenthesis g-values were not available.

FIGURES

1. Schematic diagram of flow tube. The gas inlet located 5.3 cm before the observation window was present for one experiment in lieu of the inlet right at the observation window.
2. Condensed discharge tube and firing circuit.
3. Sensitivity curve for 3/4 m Czerny-Turner monochromator with EMI-9558QA photomultiplier tube.
4. Relative steady-state population of some nickel states in the nickel carbonyl flame at 0.13 torr.
5. Relative steady-state population of some nickel states in the nickel carbonyl flame at 0.31 torr.
6. Relative steady-state population of some nickel states in the nickel carbonyl flame at 0.4 torr.
7. Relative steady-state population of some nickel states in the nickel carbonyl flame at 1.1 torr.
8. Relative steady-state population of some nickel states in the nickel carbonyl flame at 2.4 torr.
9. Relative steady-state population of some nickel states in the nickel carbonyl flame at 9.6 torr.
10. Relative steady-state population of some nickel states in the nickel carbonyl flame at 14.7 torr.
11. Relative steady-state population of some nickel states in the nickel carbonyl flame at 27 torr.
12. Effective temperature of the relative steady-state population distribution as a function of pressure.
13. Rate of formation of some nickel states at 0.13 torr.
14. Rate of formation of some nickel states at 2.4 torr. Calculated values are shown as squares.
15. Rates of formation of some nickel states at 27 torr. Calculated values for a pressure of 40 torr are represented as squares.
16. Quenching of the long-lived 30980 cm^{-1} state of nickel.

(a)



(b)

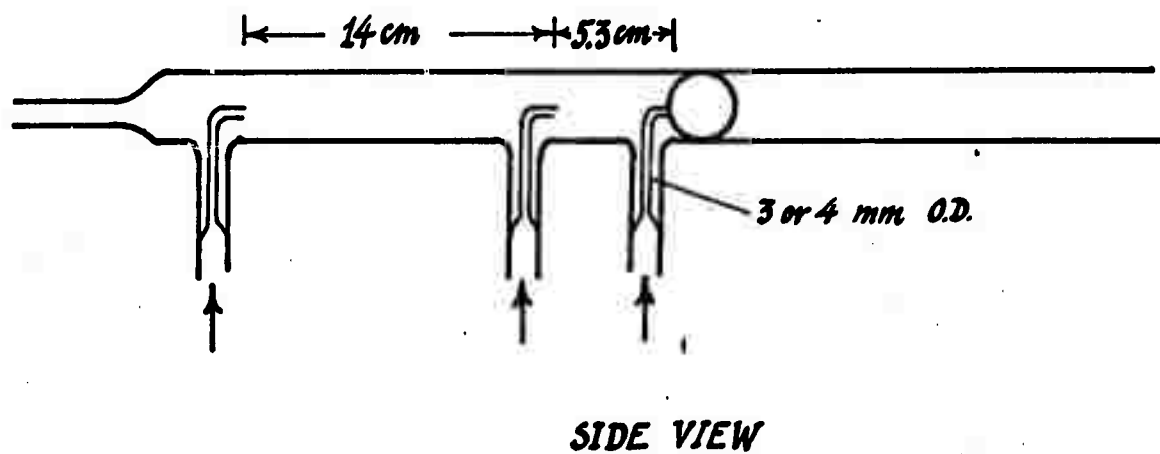
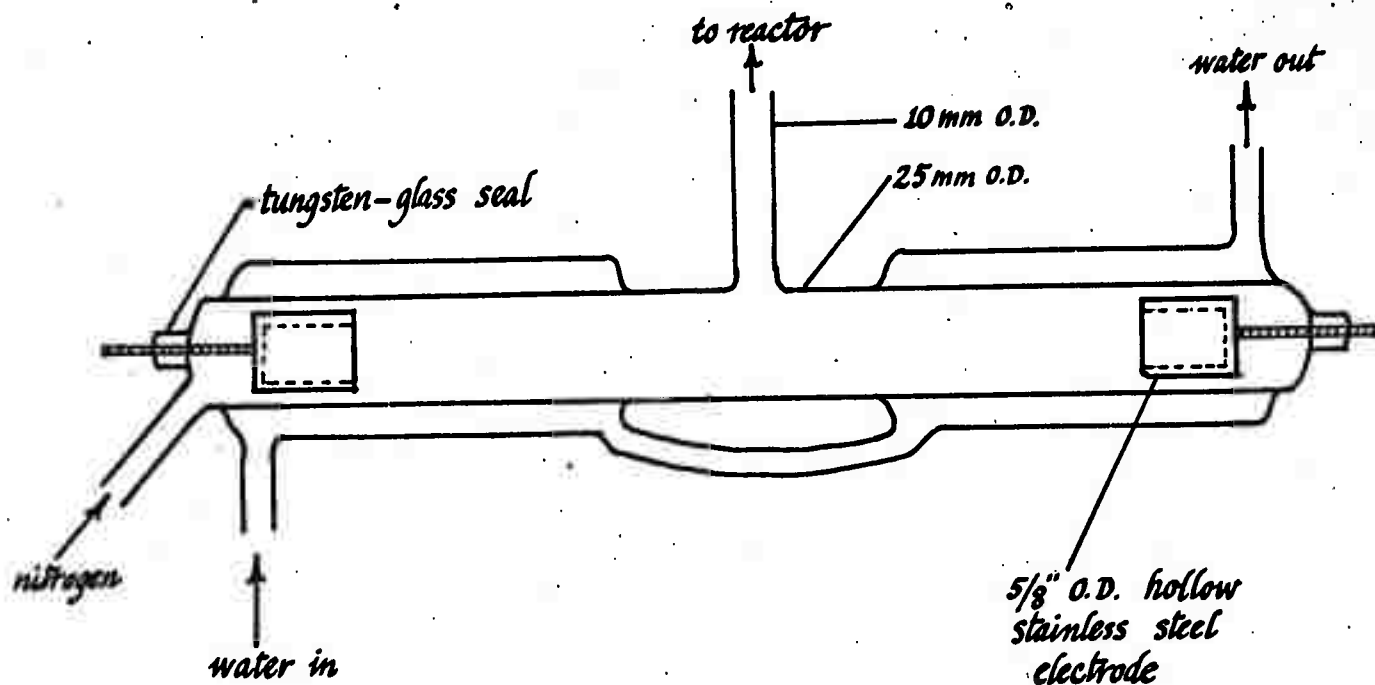
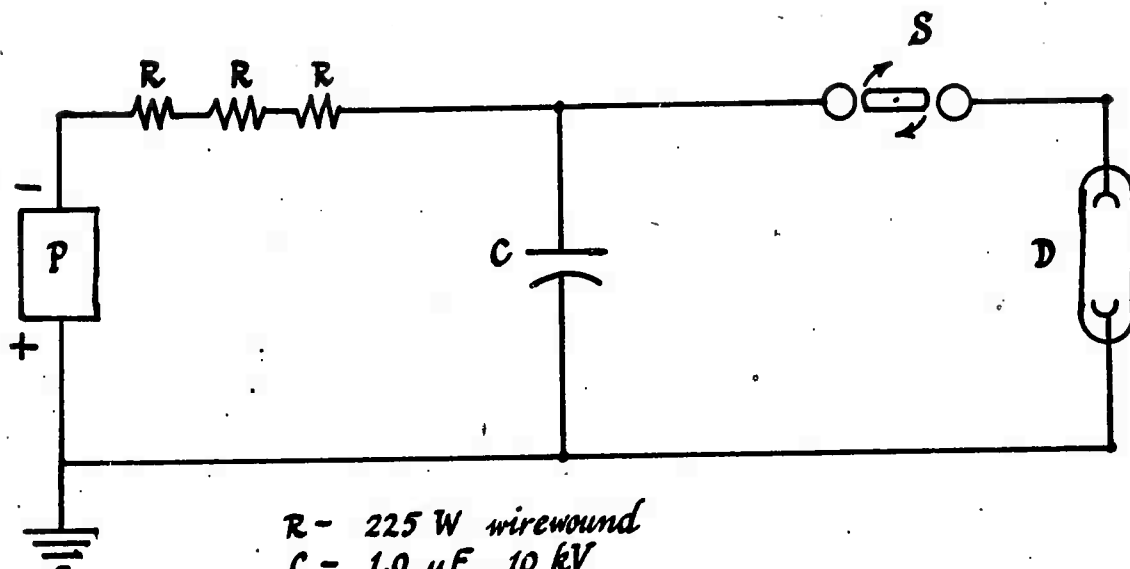


FIGURE 1



(a)



R - 225 W wirewound
 C - 1.0 μ F, 10 kV
 P - D.C. power supply
 S - rotating spark gap
 D - pulsed discharge tube

(b)

FIGURE 2

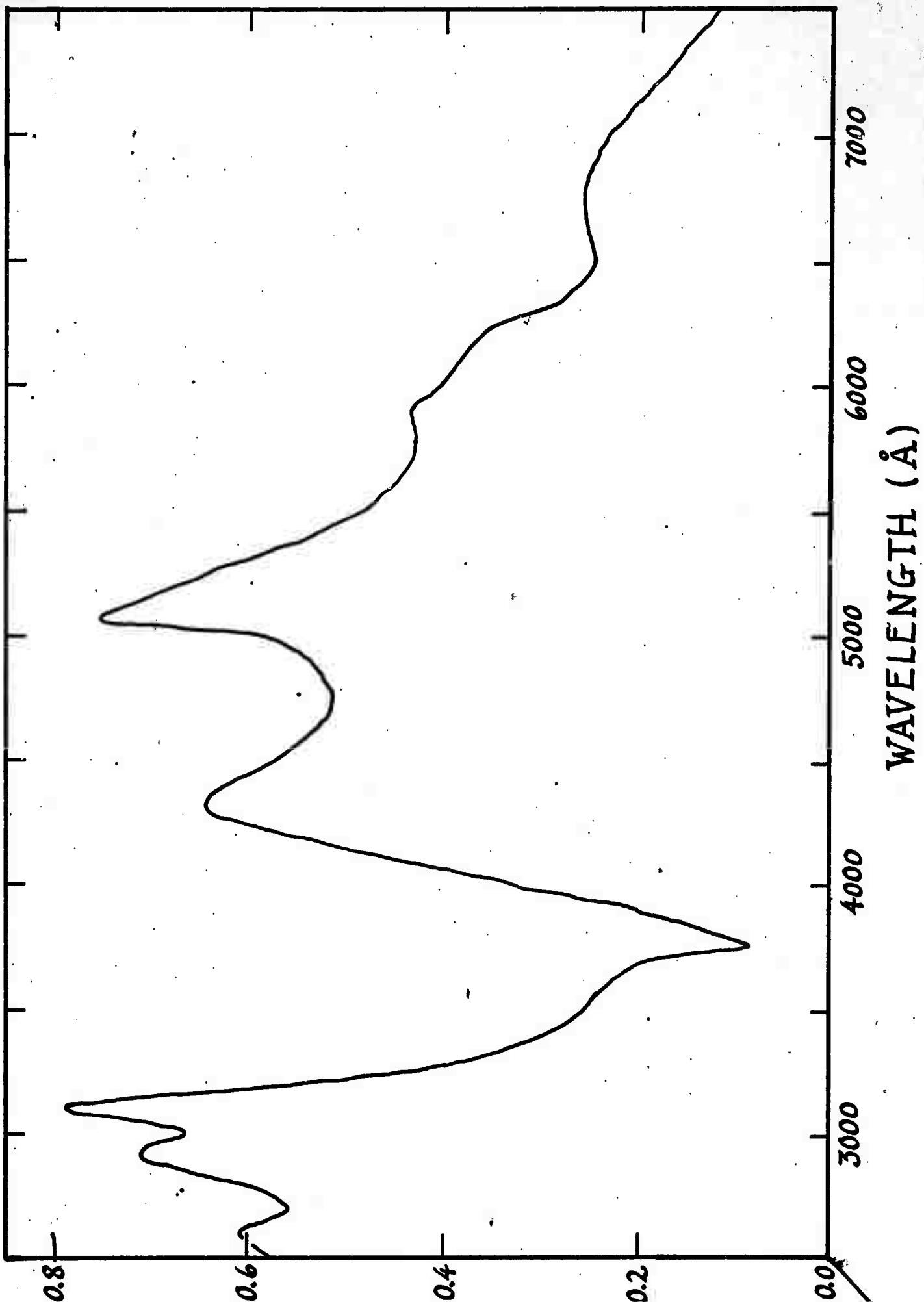


FIGURE 3

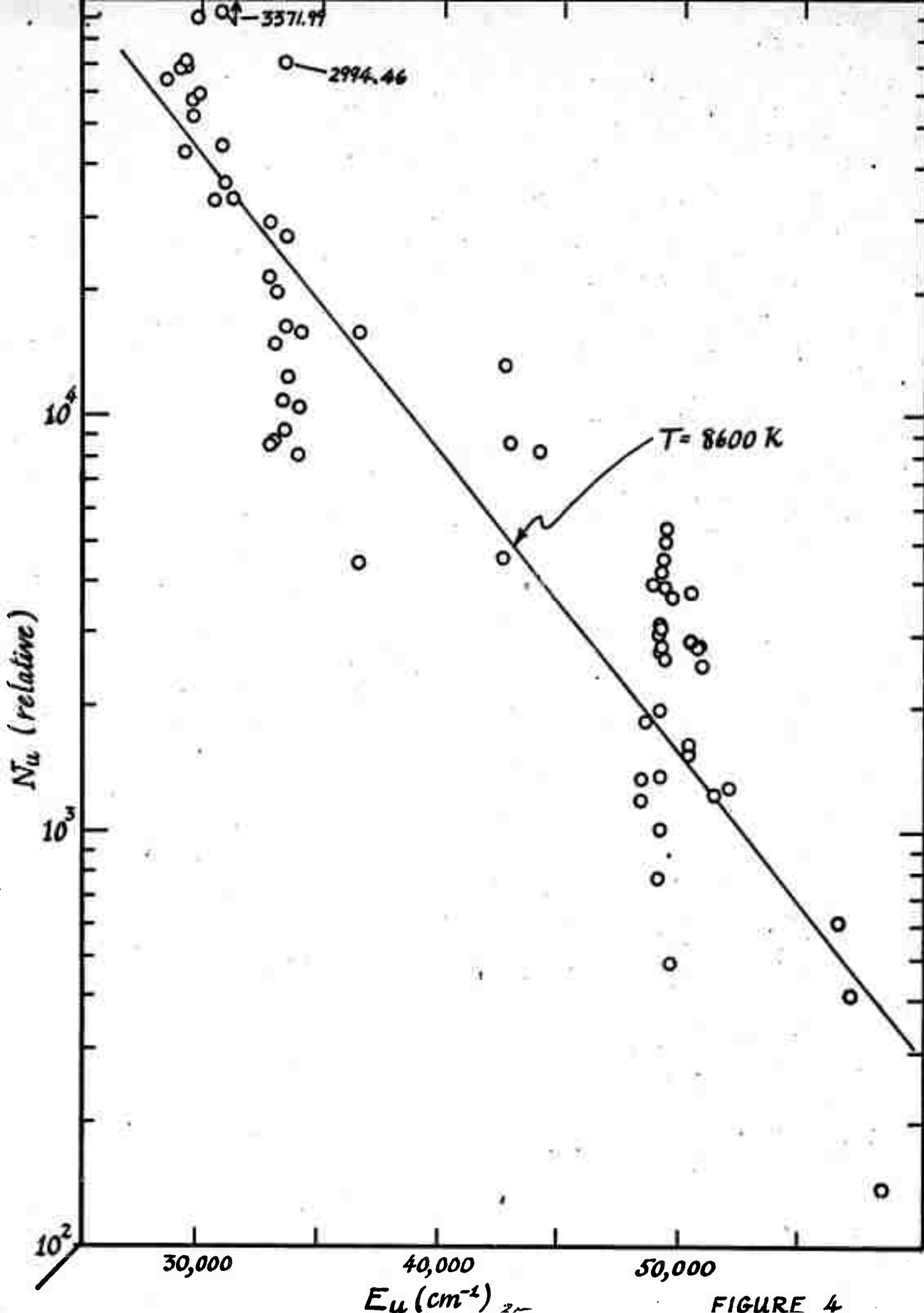
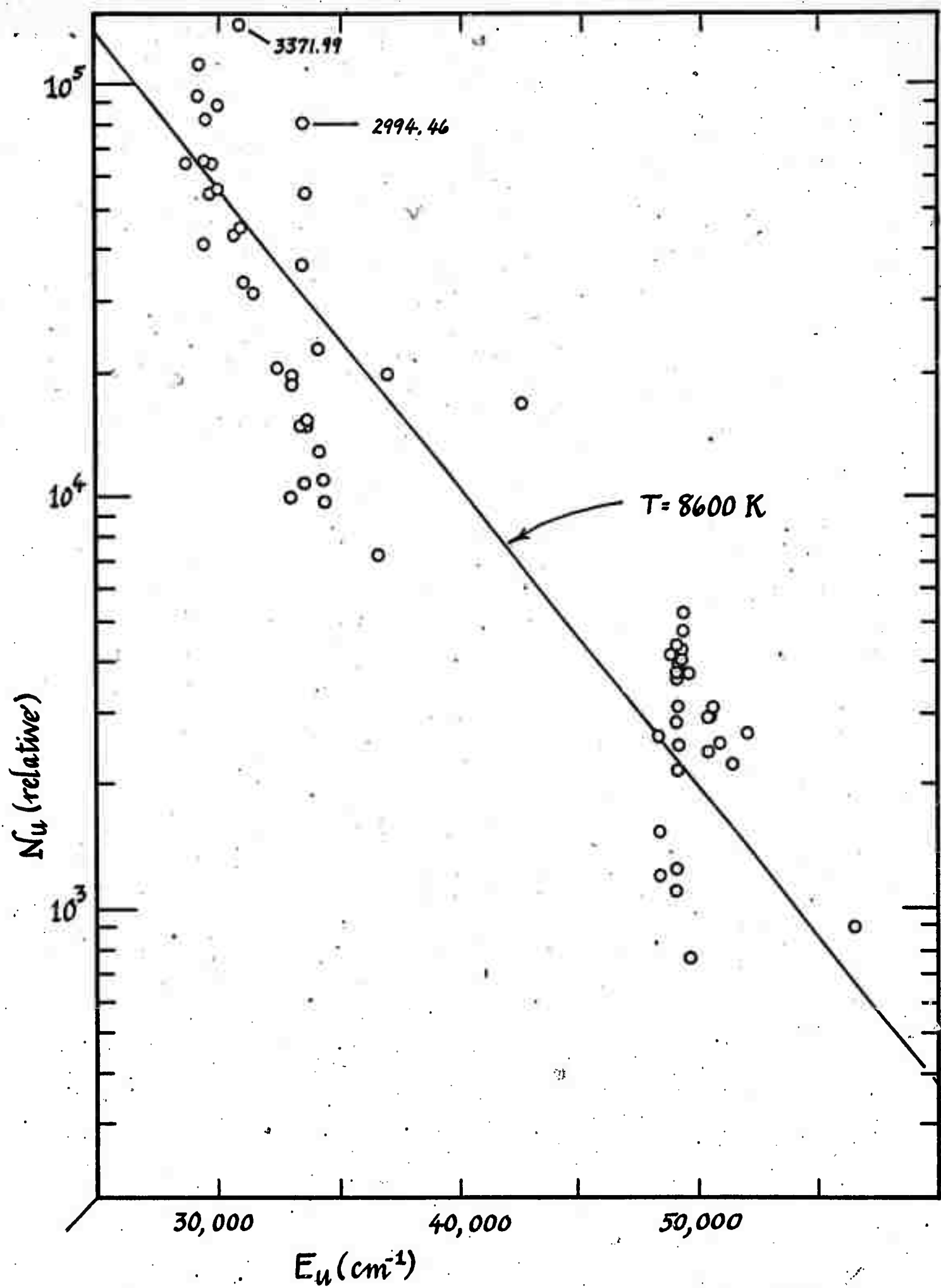


FIGURE 4



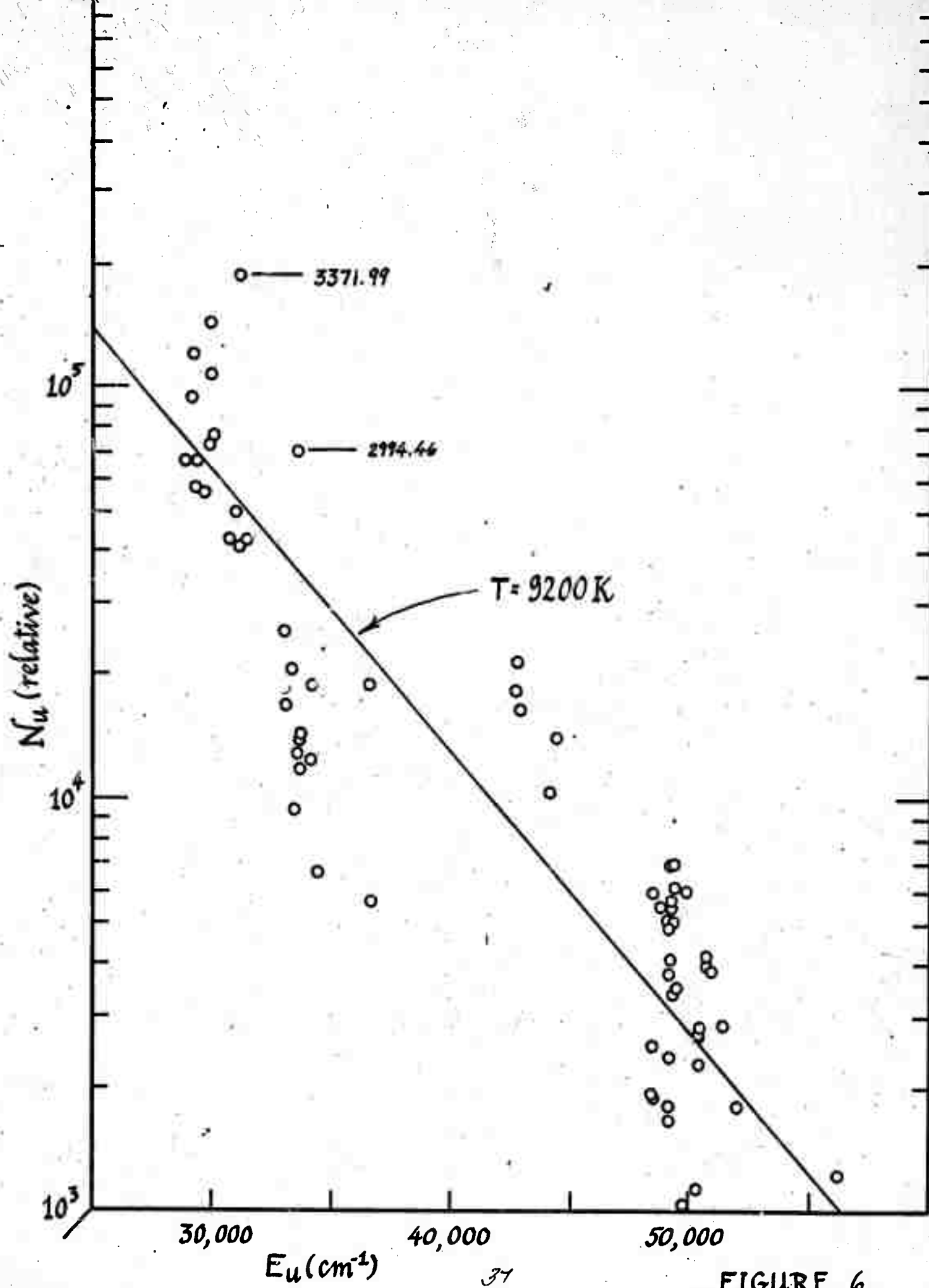


FIGURE 6

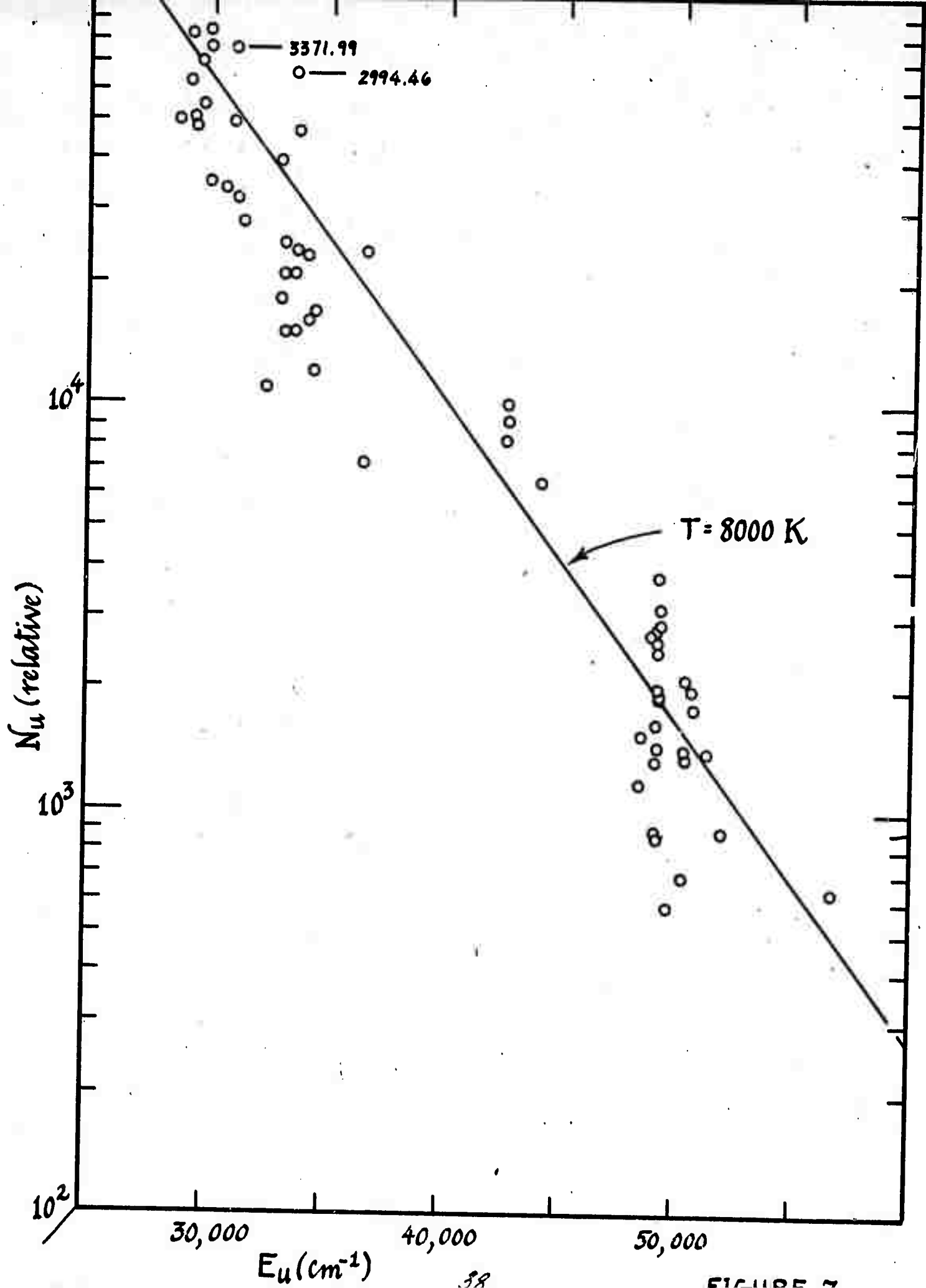
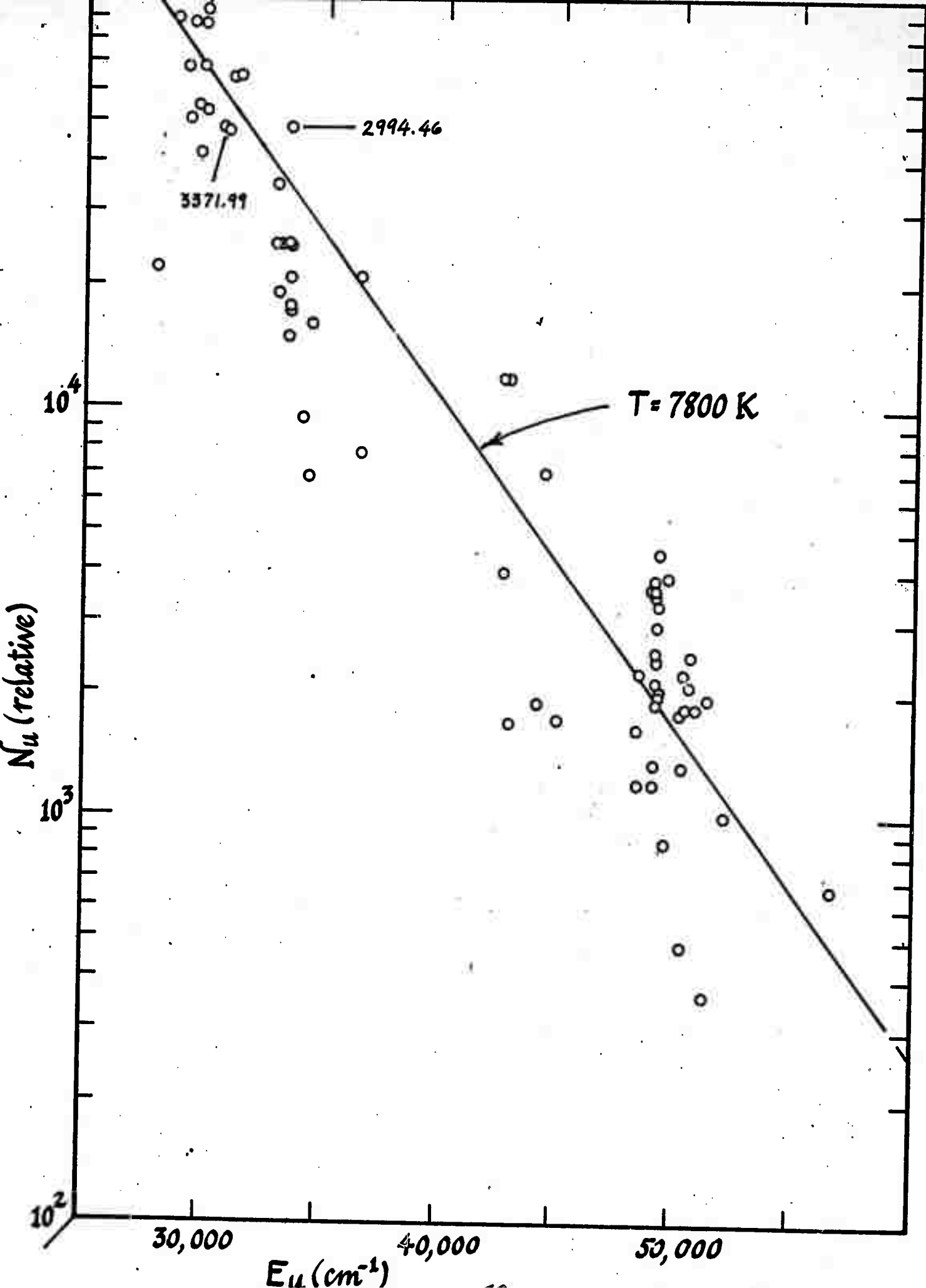
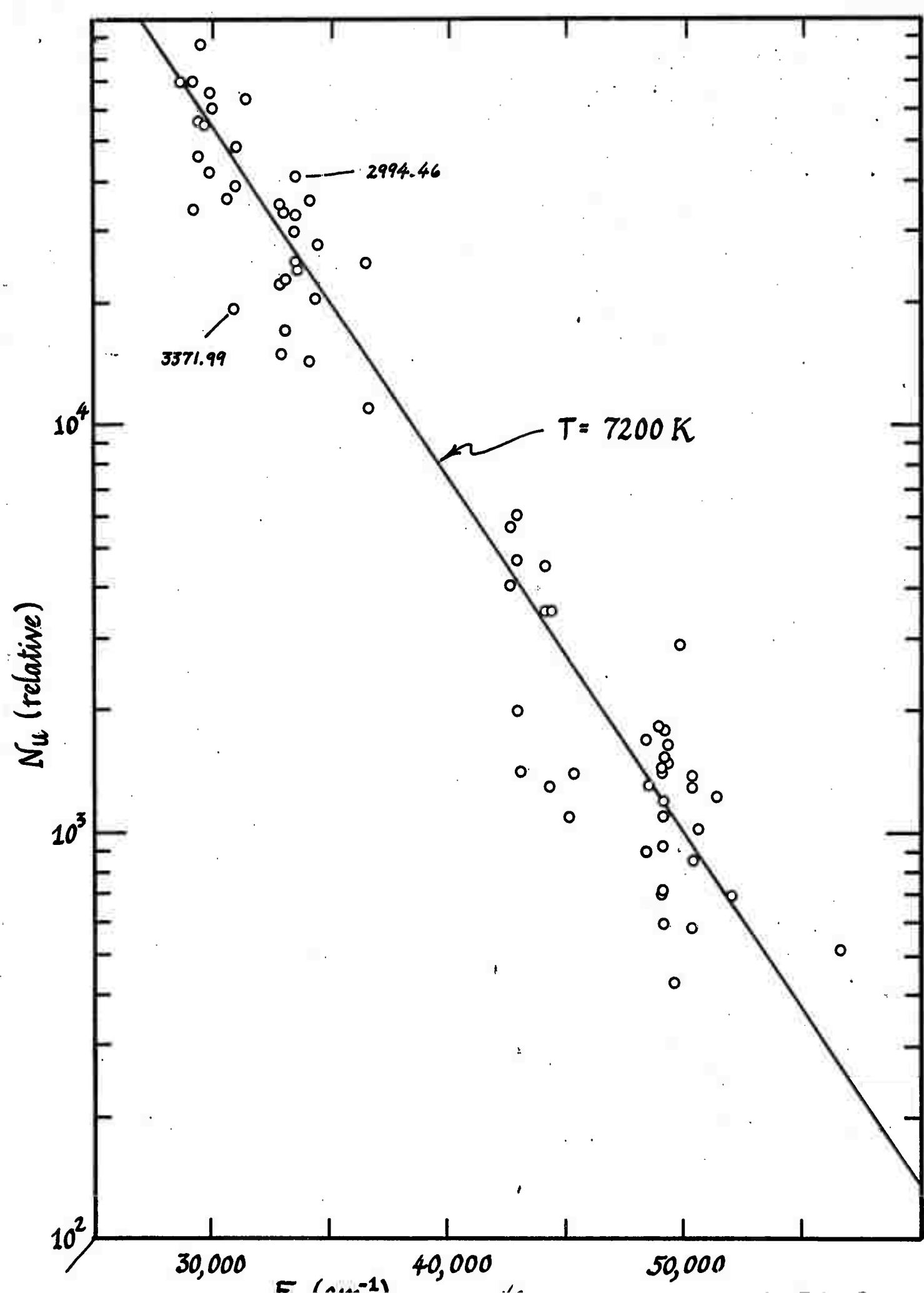
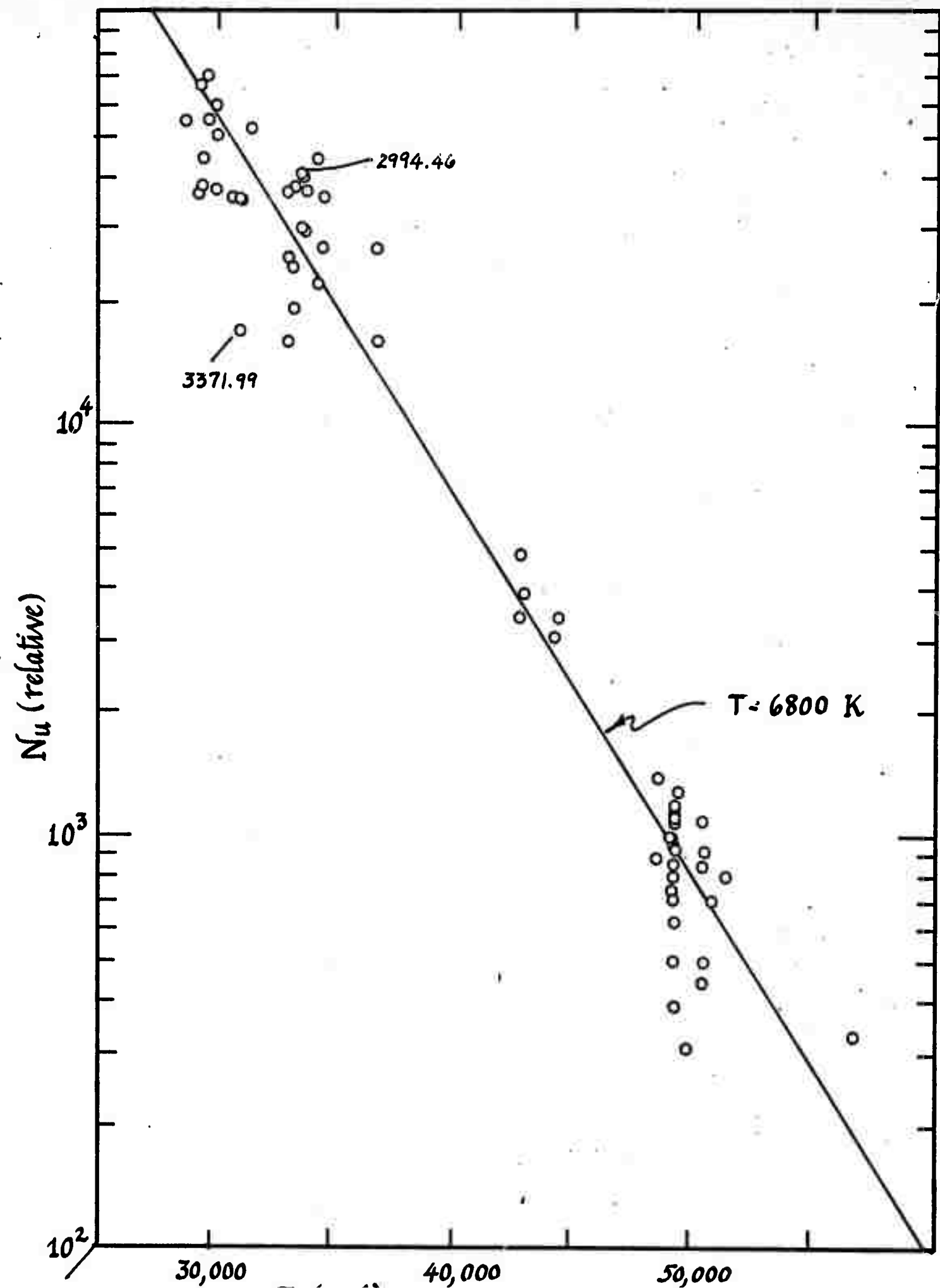


FIGURE 7







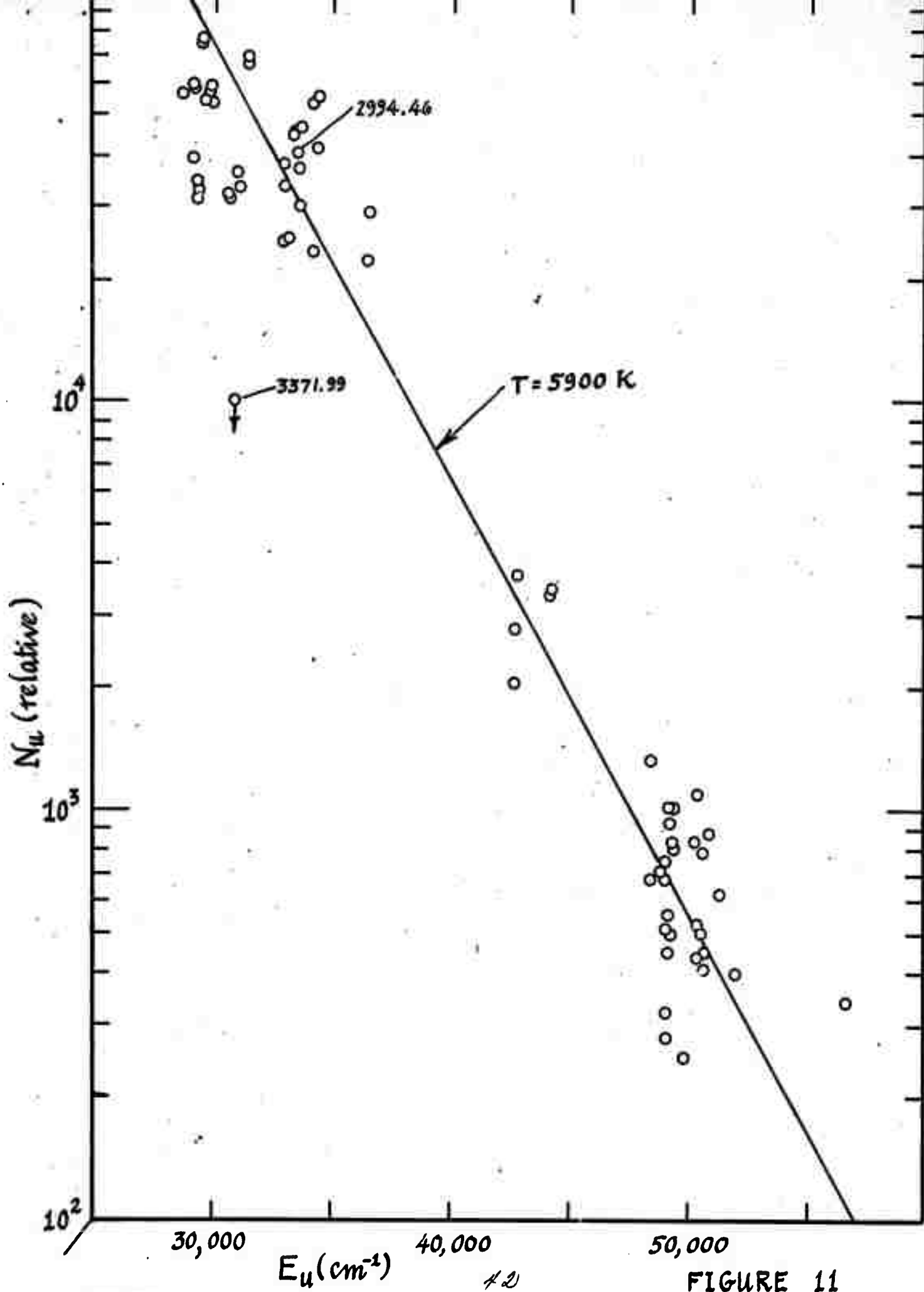


FIGURE 11

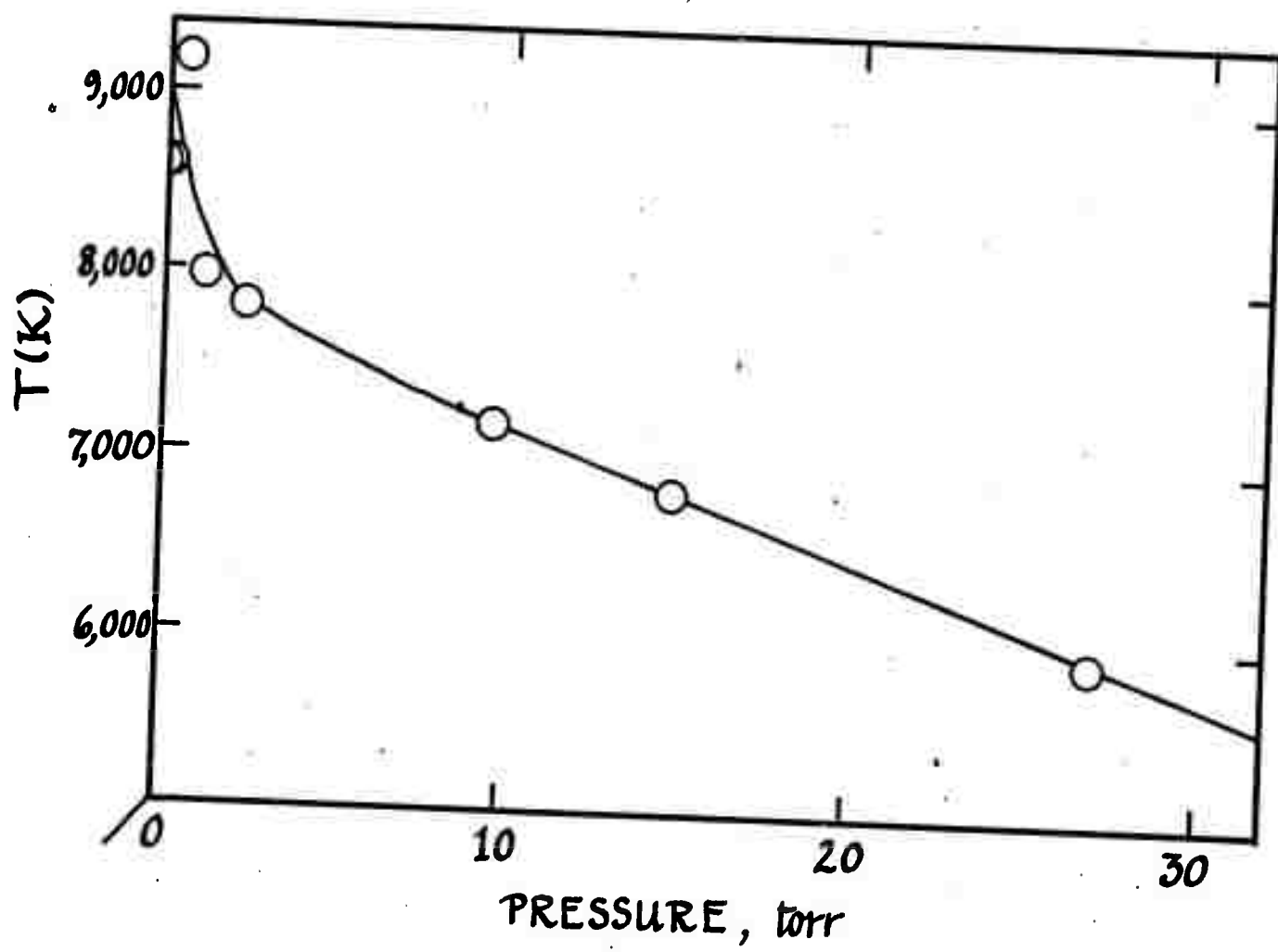


FIGURE 12

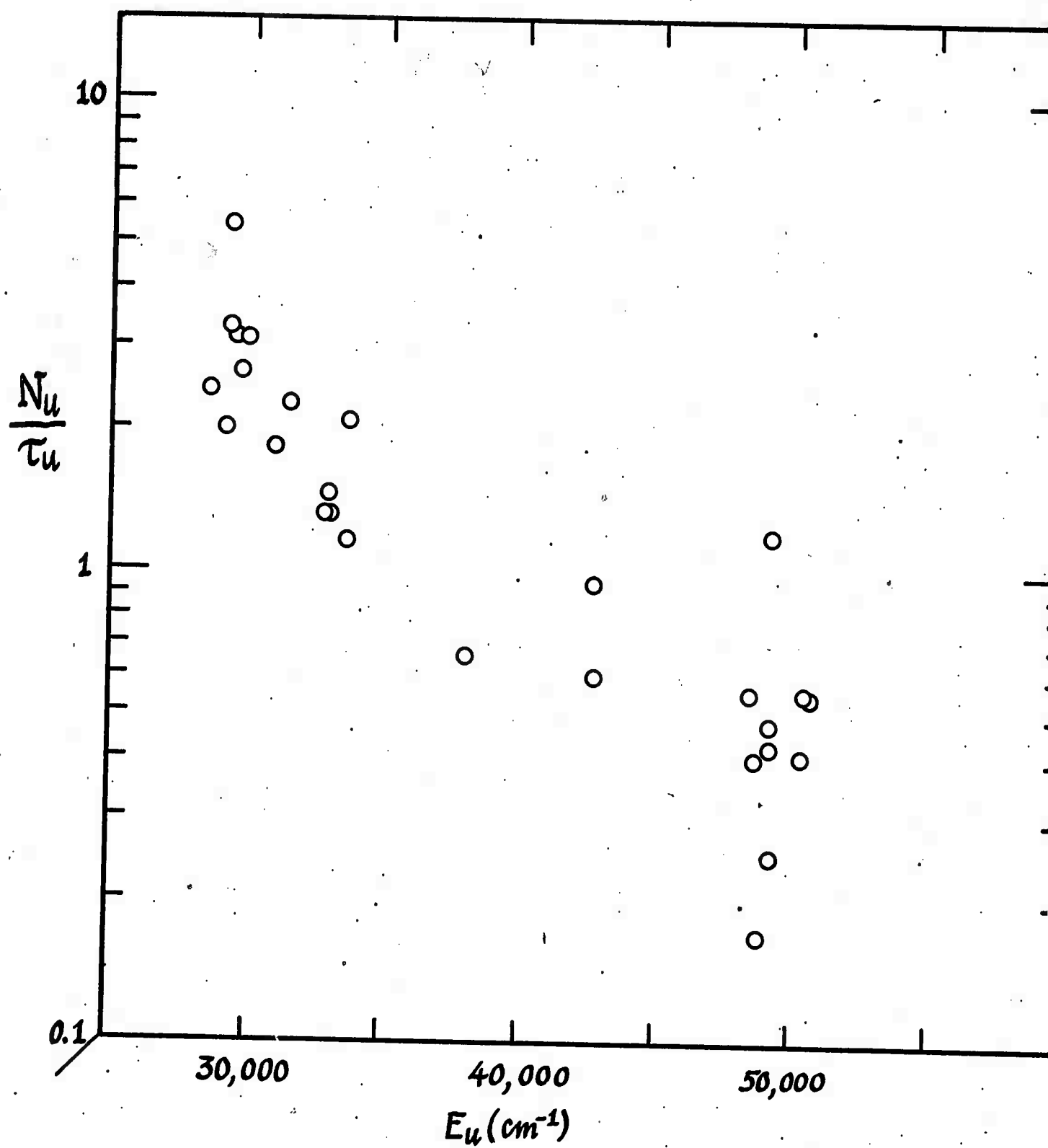
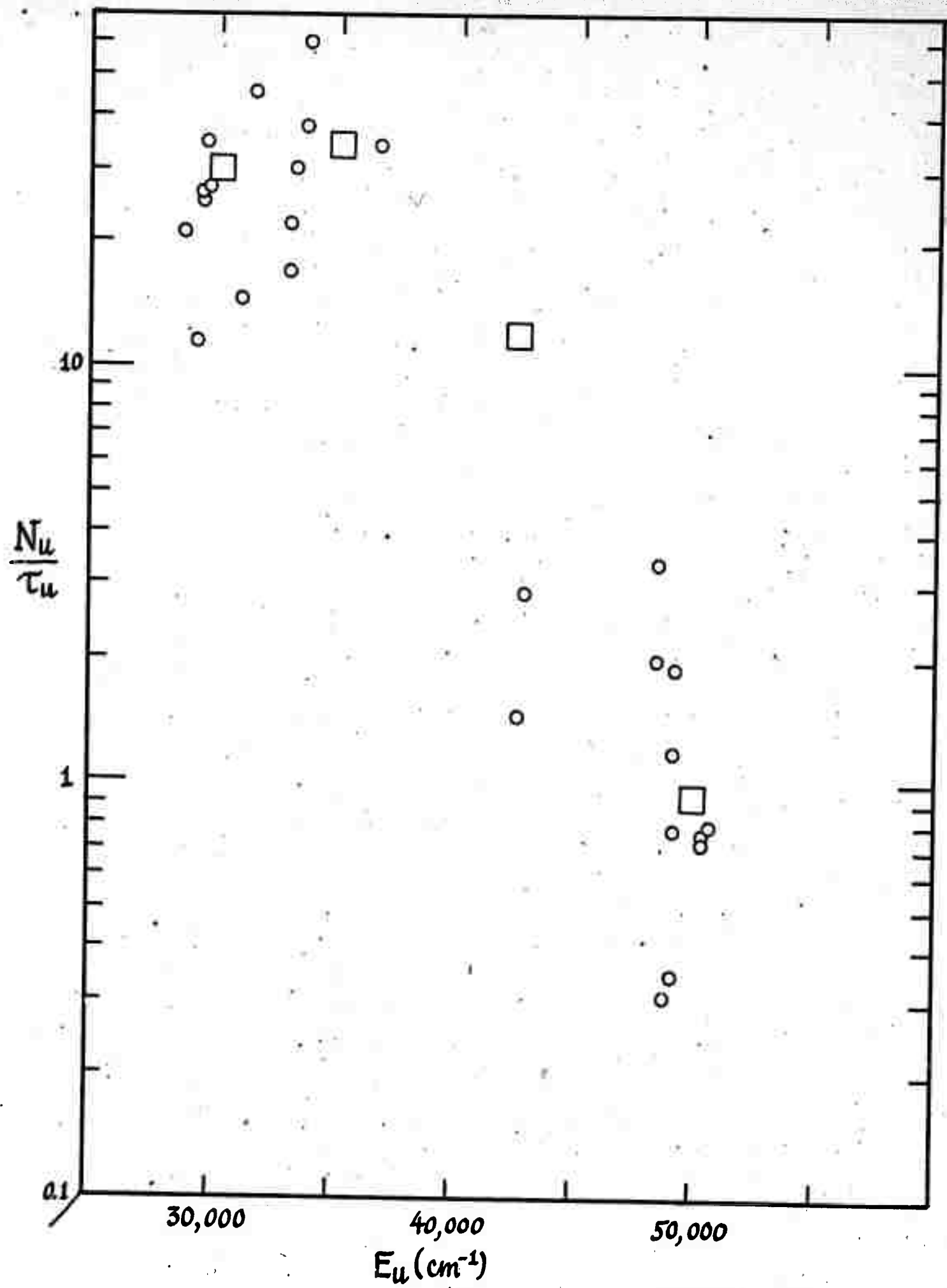


FIGURE 13



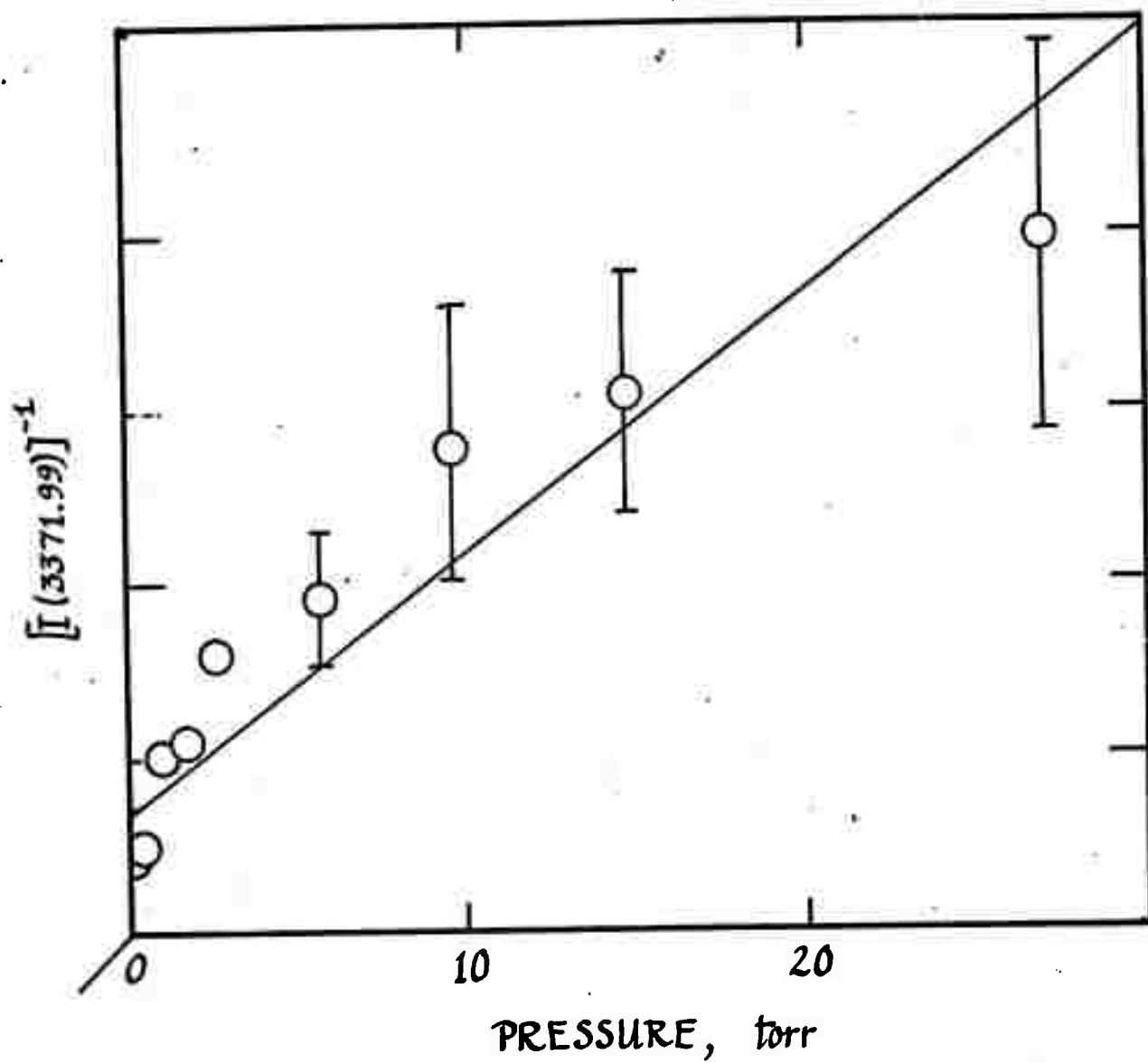


FIGURE 16